



Building 911C  
P.O. Box 5000  
Upton, NY 11973-5000  
Phone 631 344-4531  
Fax 631 344-5954  
hershcovitch@bnl.gov

*DATE:* May 25, 2007

## Memo

*TO:* RHIC E-Coolers

*FROM:* Ady Hershcovitch

*SUBJECT:* **Minutes of the May 25, 2007 Meeting**

Present: Natalia Abreu, Michael Blaskiewicz, Mike Brennan, Alexei Fedotov, Wolfram Fischer, Harald Hahn, Ady Hershcovitch, Derek Lowenstein, William Mackay, Eduard Pozdeyev, Thomas Roser, Anatoly Sidorin (JINR Dubna Russia), Dejan Trbojevic, Gang Wang, Alexander Zaltsman.

Topic discussed: Stochastic Cooling

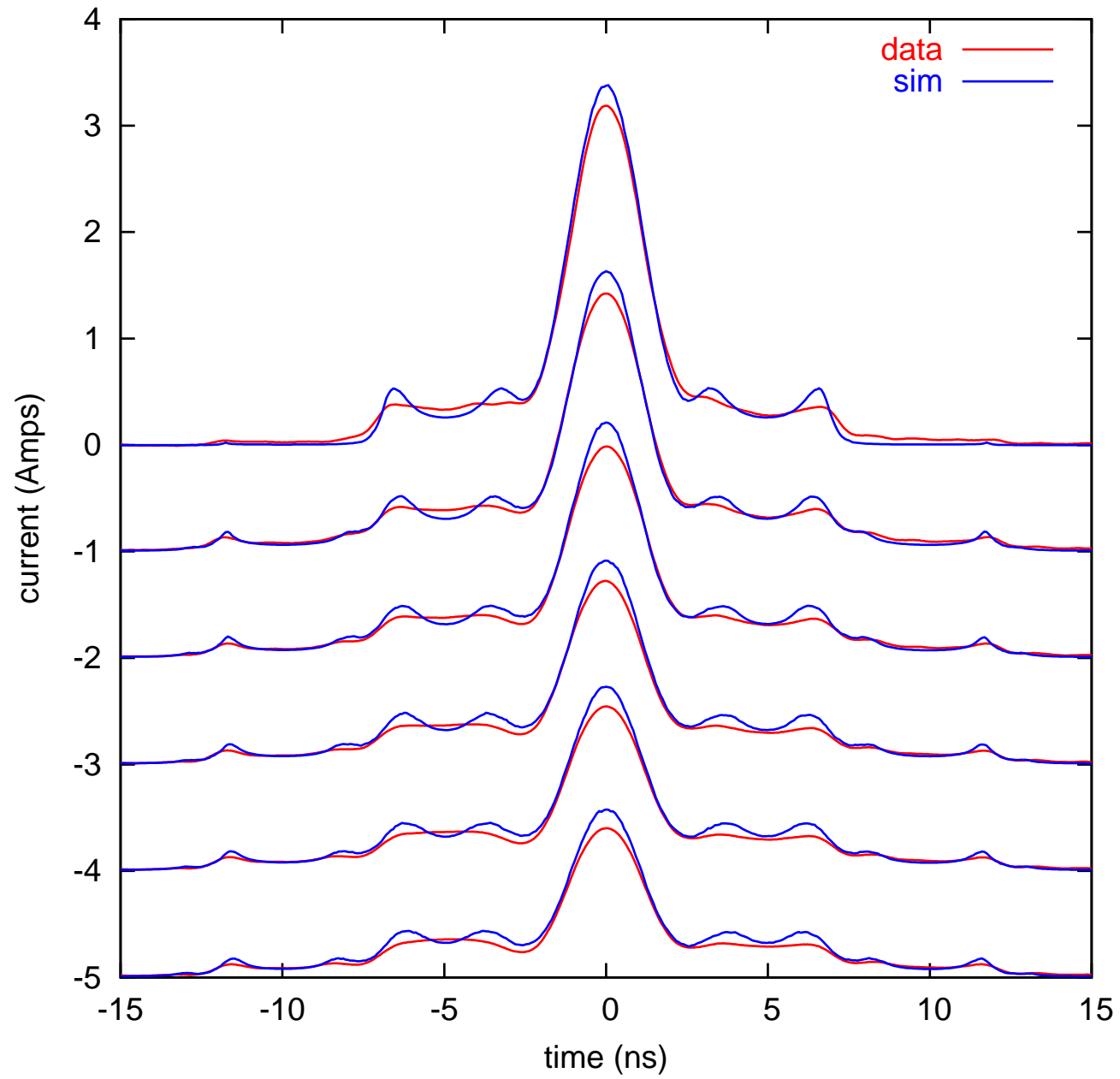
**Stochastic Cooling:** during the meeting Mike Blaskiewicz presented theory, simulations and experimental results of stochastic cooling in RHIC. Mike started with a short review of the theory behind stochastic cooling. Basically, transverse stochastic cooling is a transverse wide-band damper comprising of pickups and kickers. First Mike described cooling theory without intrabeam scattering (IBS). Next came a description of Piwinski's theory of intrabeam scattering. The theory assumes the bunch is Gaussian in all three dimensions, which is a poor approximation to rebucketed beam in RHIC. Mike modified the theory by assuming that the local intrabeam scattering rate is proportional to the local line density of the particles. The rms values of momentum spread and betatron action are taken as constant along the bunch and the IBS kicks are normalized so that the rms parameters of a Gaussian bunch will evolve as they do in Piwinski's original equations. Additionally, since the cooling rate is inversely proportional to the number of particles, a simulation with  $10^5$  particles can be scaled to represent the behavior of a bunch with  $10^9$  ions. Mike introduced the same scaling into the IBS rates and stated that an internal consistency check had verified that everything scaled as expected. The simulations track only the x transverse dimension and assume that the emittance in the y transverse dimension is the same. There is one free parameter, the ratio of the x emittance growth rate to the value calculated using Piwinski's formulae. This ratio,  $f_{\text{bsx}}$ , was set to either 0.5, or 1.

Graphs below show the results for several simulations. Plots 1 and 2 compare data and simulations for the effect of intrabeam scattering alone. Plot 3 compares the evolution of the transverse emittance in the two cases. Plot 4 shows the transverse emittance when the longitudinal cooling is on and strong. Plots 5 and 6 compare data and simulations for longitudinal cooling on. Plots 7 and 8 compare data and simulation with the cooling rate at 2/3 optimal. Plots 9 and 10 show the effect of fairly strong transverse cooling on the longitudinal profile. It is likely that  $f_{\text{bsx}}=0.5$  is more realistic than  $f_{\text{bsx}}=1$ . Comparing the

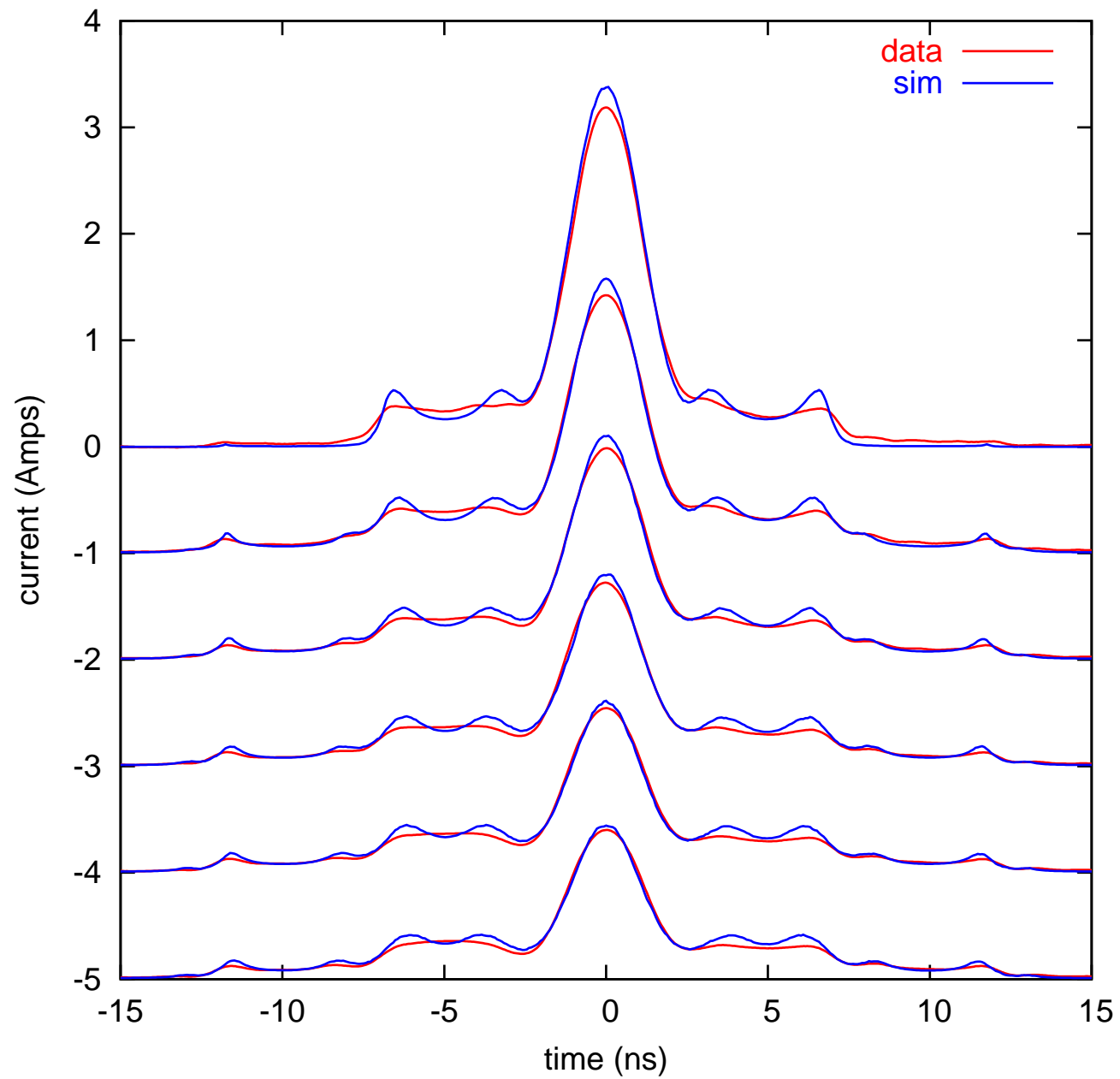
bottom trace on plot 6 and the top trace on plot 10 one can see that reducing the transverse emittance increases the longitudinal IBS rate. Plots 11 and 12 show the longitudinal mountain ranges for a smaller transverse cooling gain. There is less leakage into the satellites. Plots 13 and 14 show the evolution of the transverse emittance with cooling.

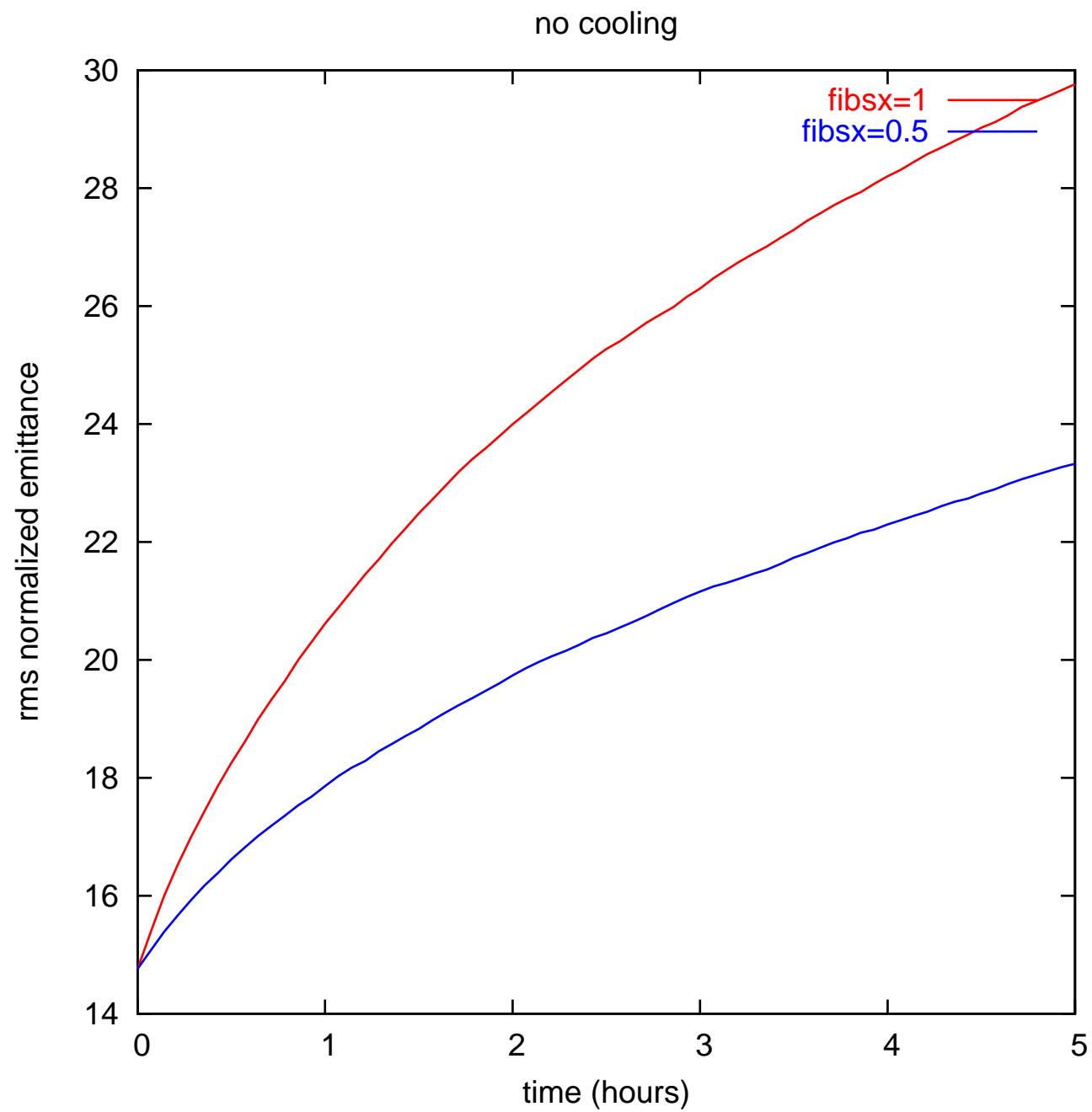
The meeting ended with a discussion regarding schemes to move ions from satellites back into the buckets.

fibsx=1, no cooling, fill 8794

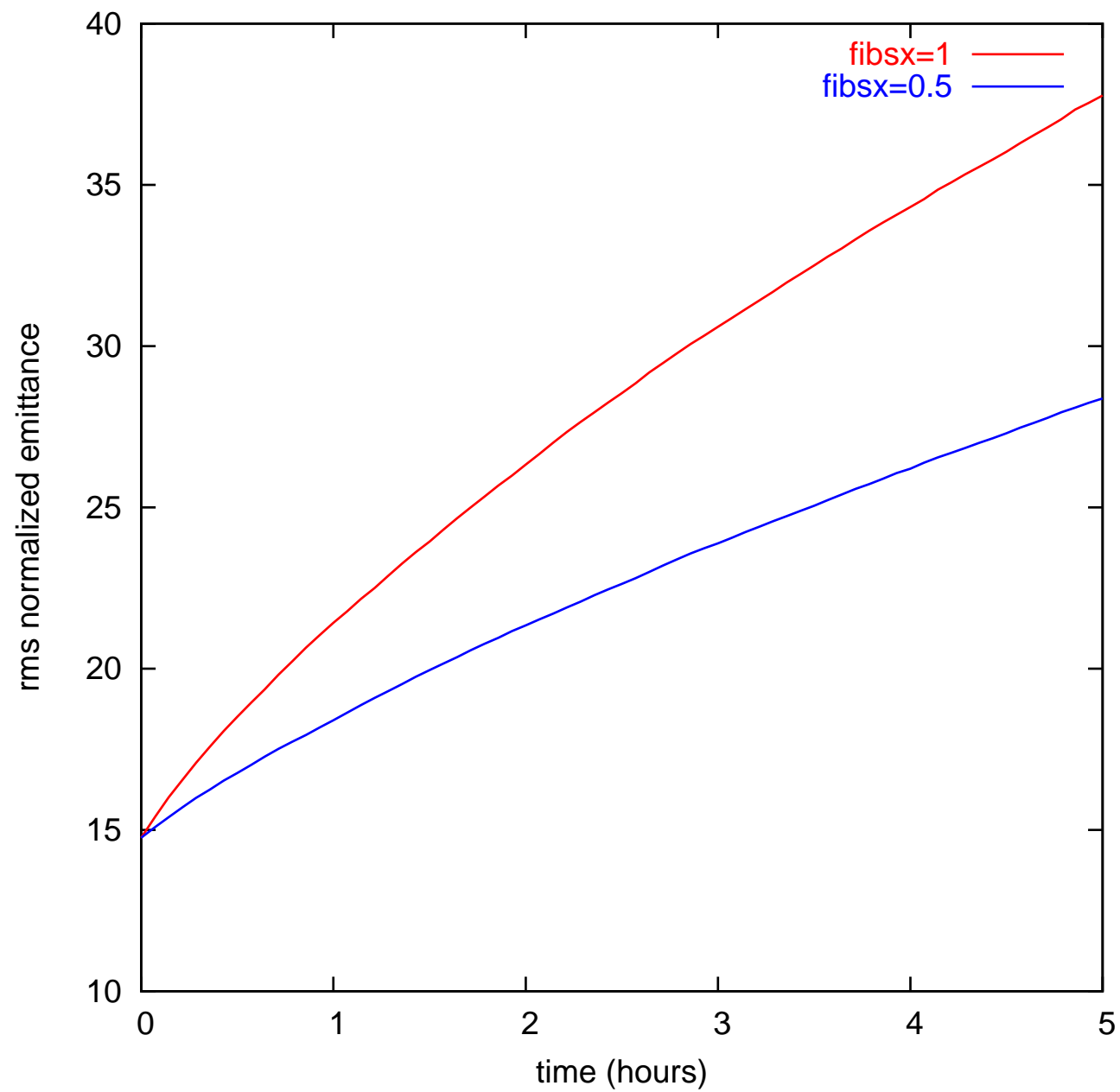


fibsx=0.5, no cooling, fill 8794

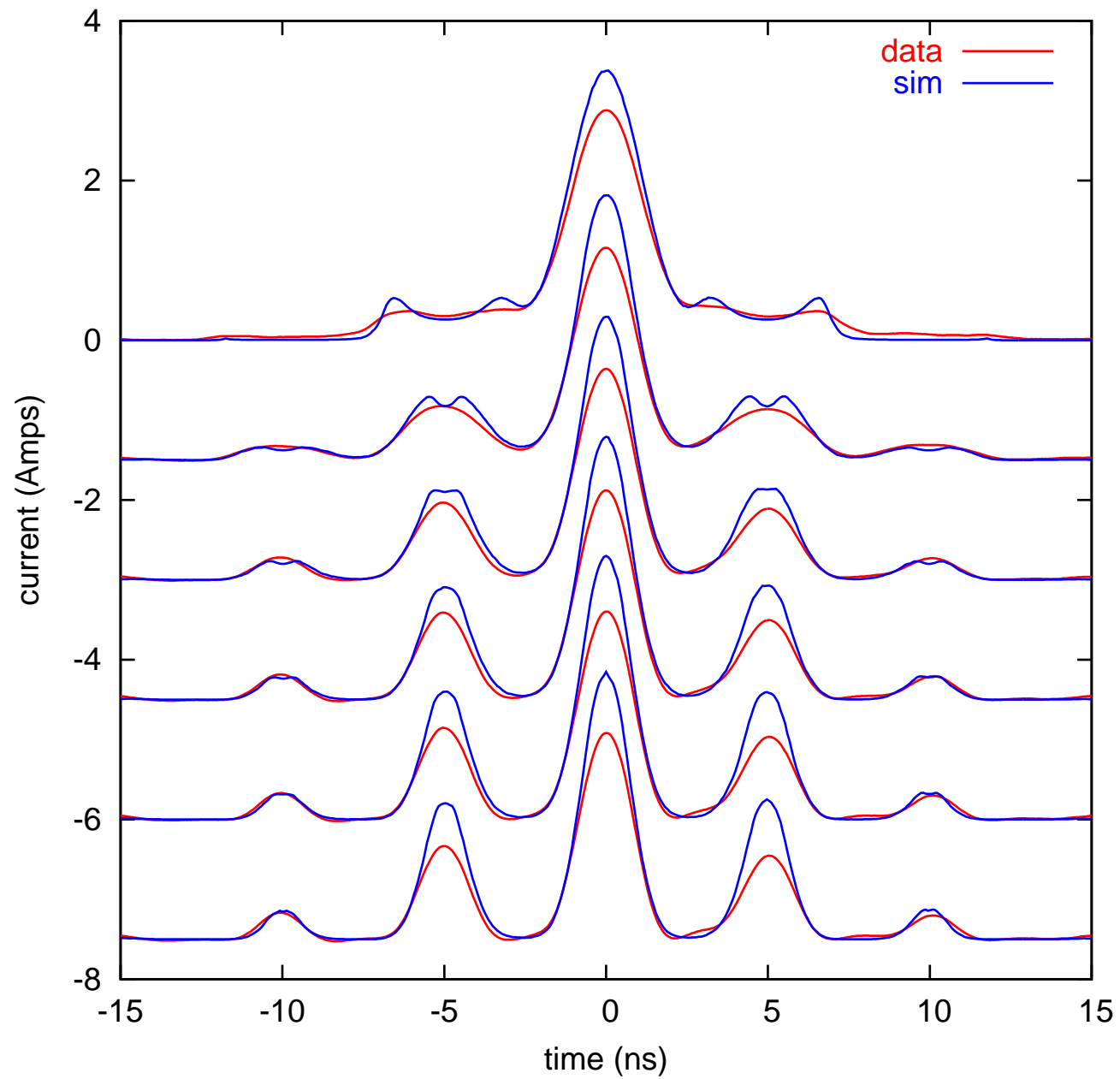




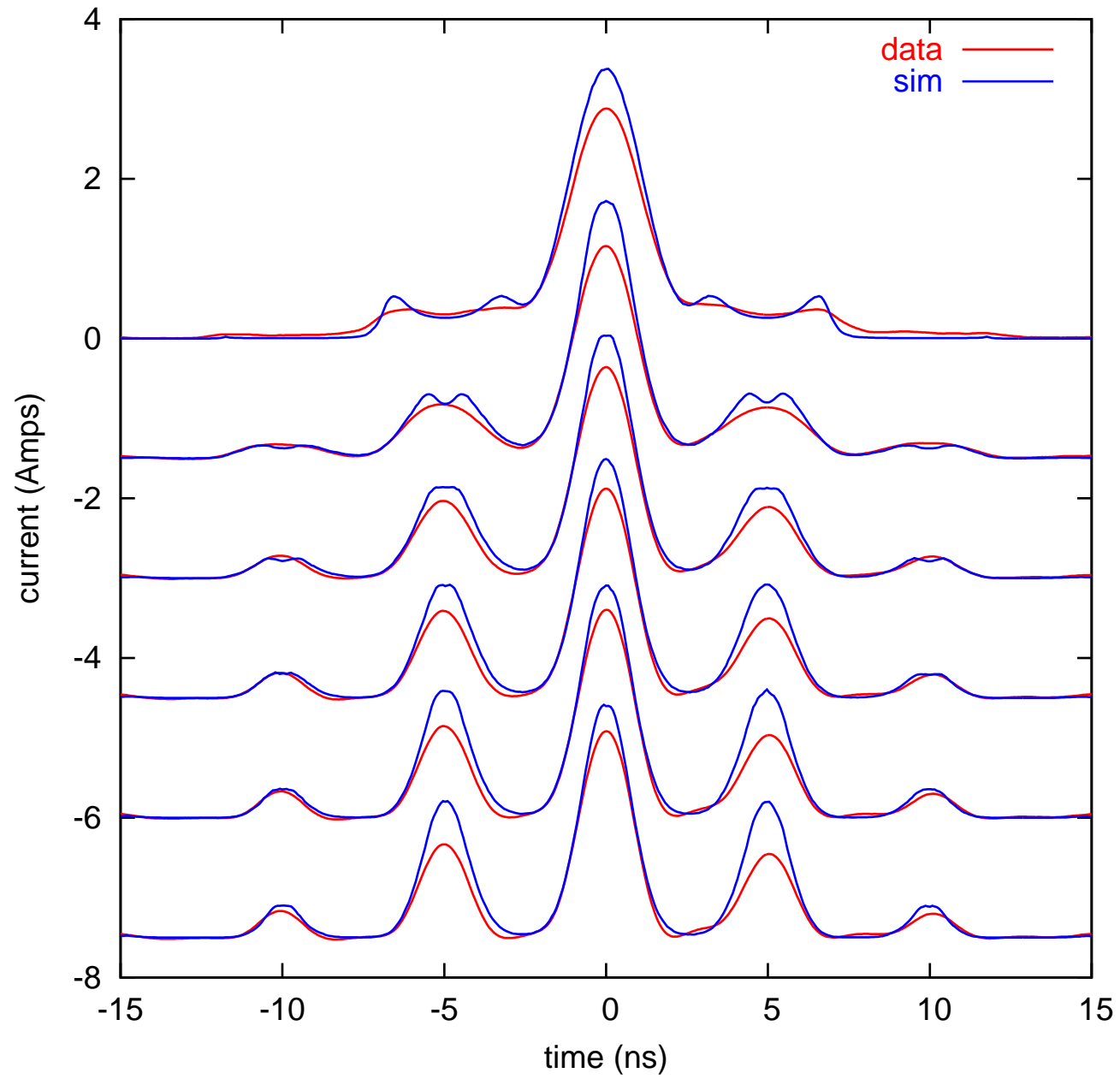
$Z_s=3.e7$



fibsx=1, Zs=3.e7, fill 8794

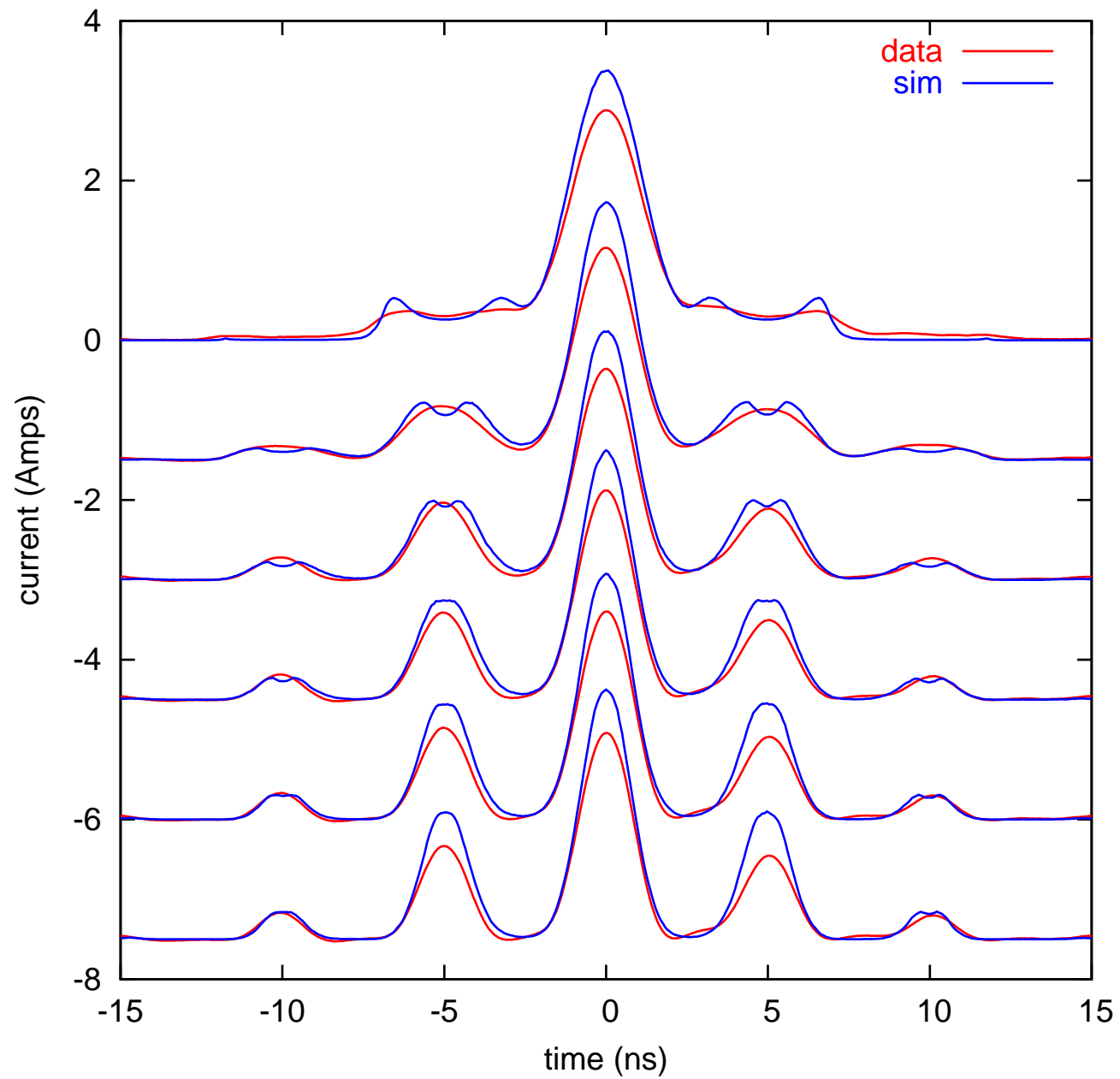


fibsx=0.5, Zs=3.e7, fill 8794

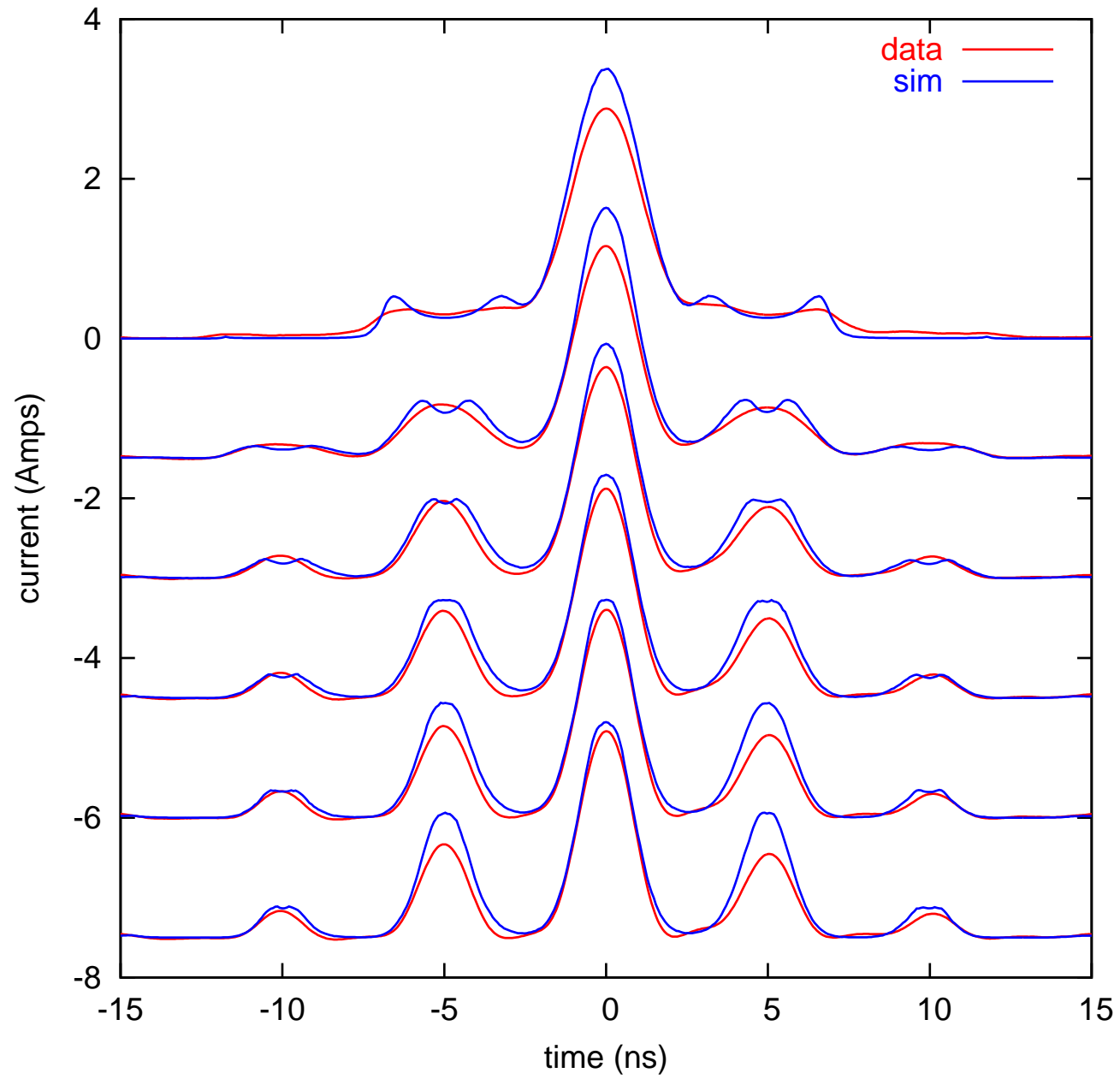


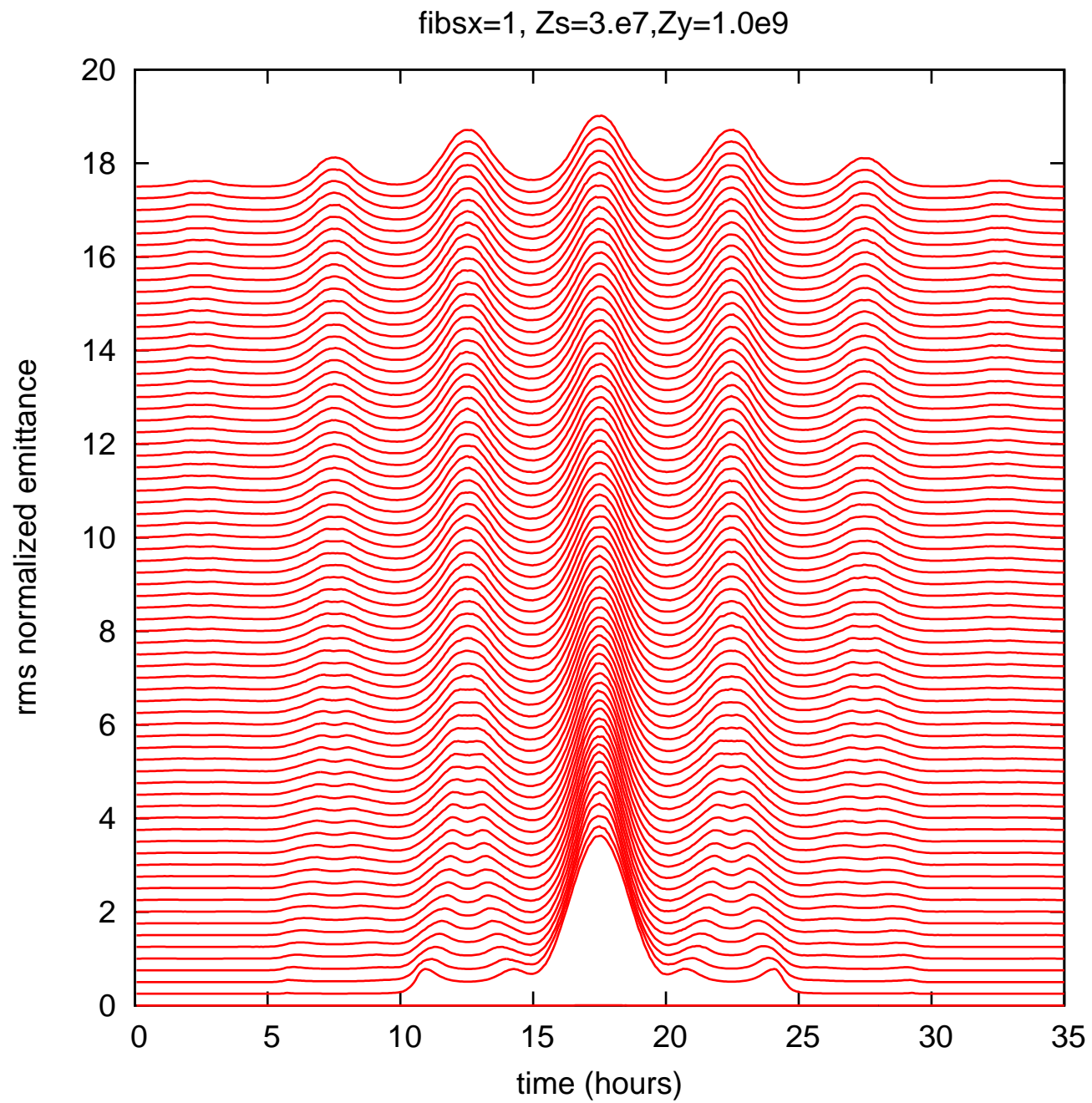


fibsx=1, Zs=2.e7, fill 8794

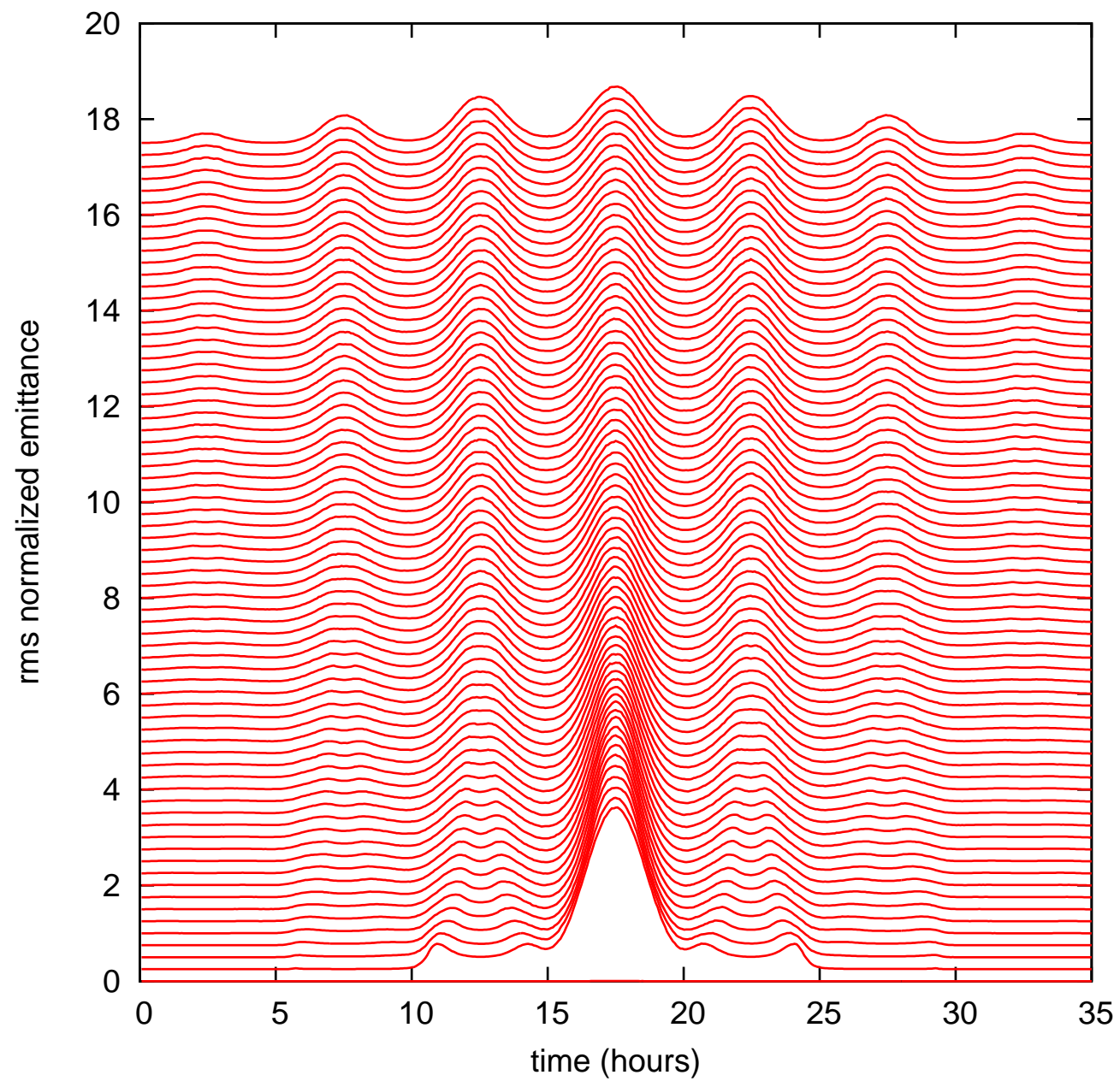


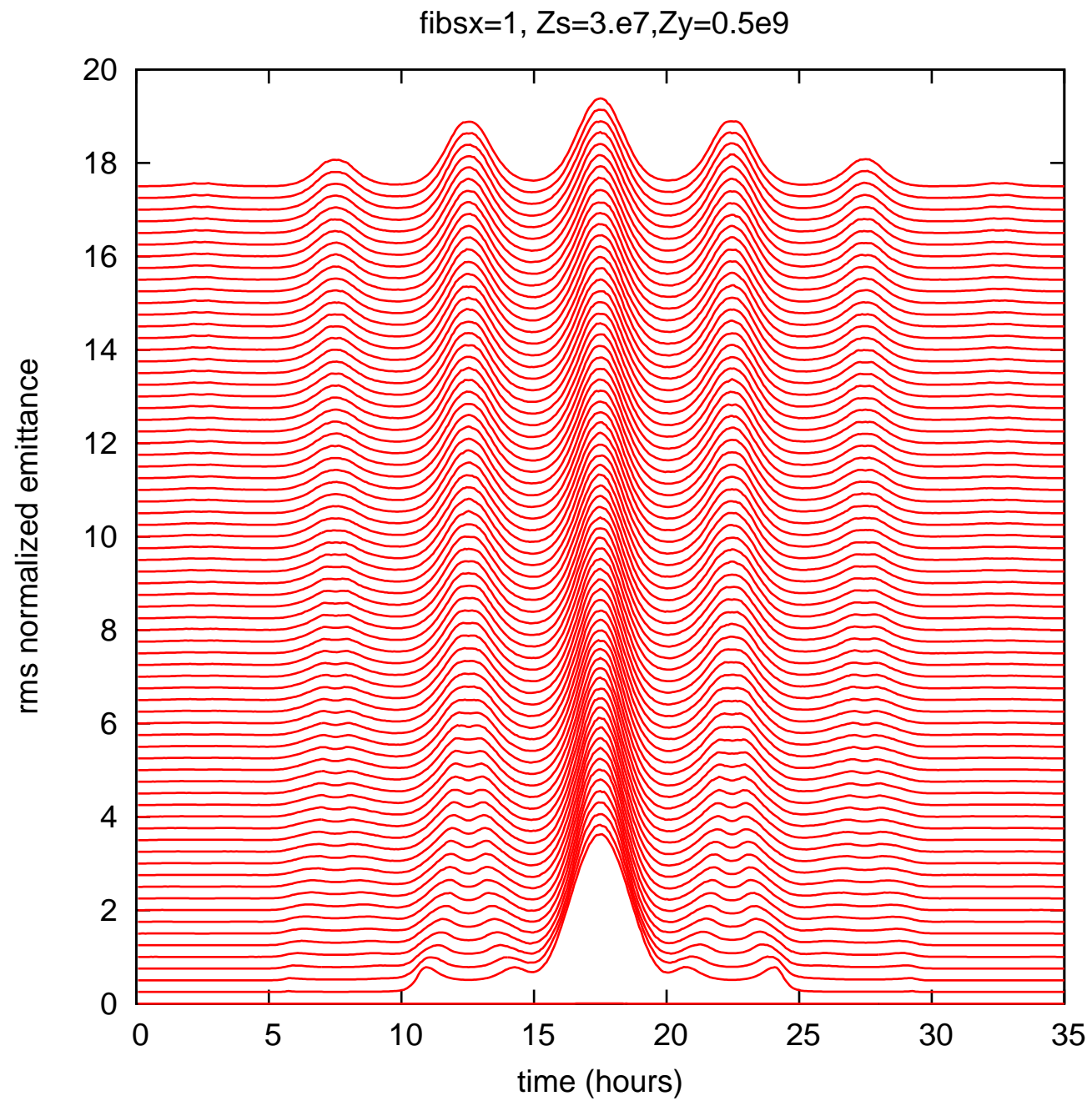
fibsx=0.5, Zs=2.e7, fill 8794



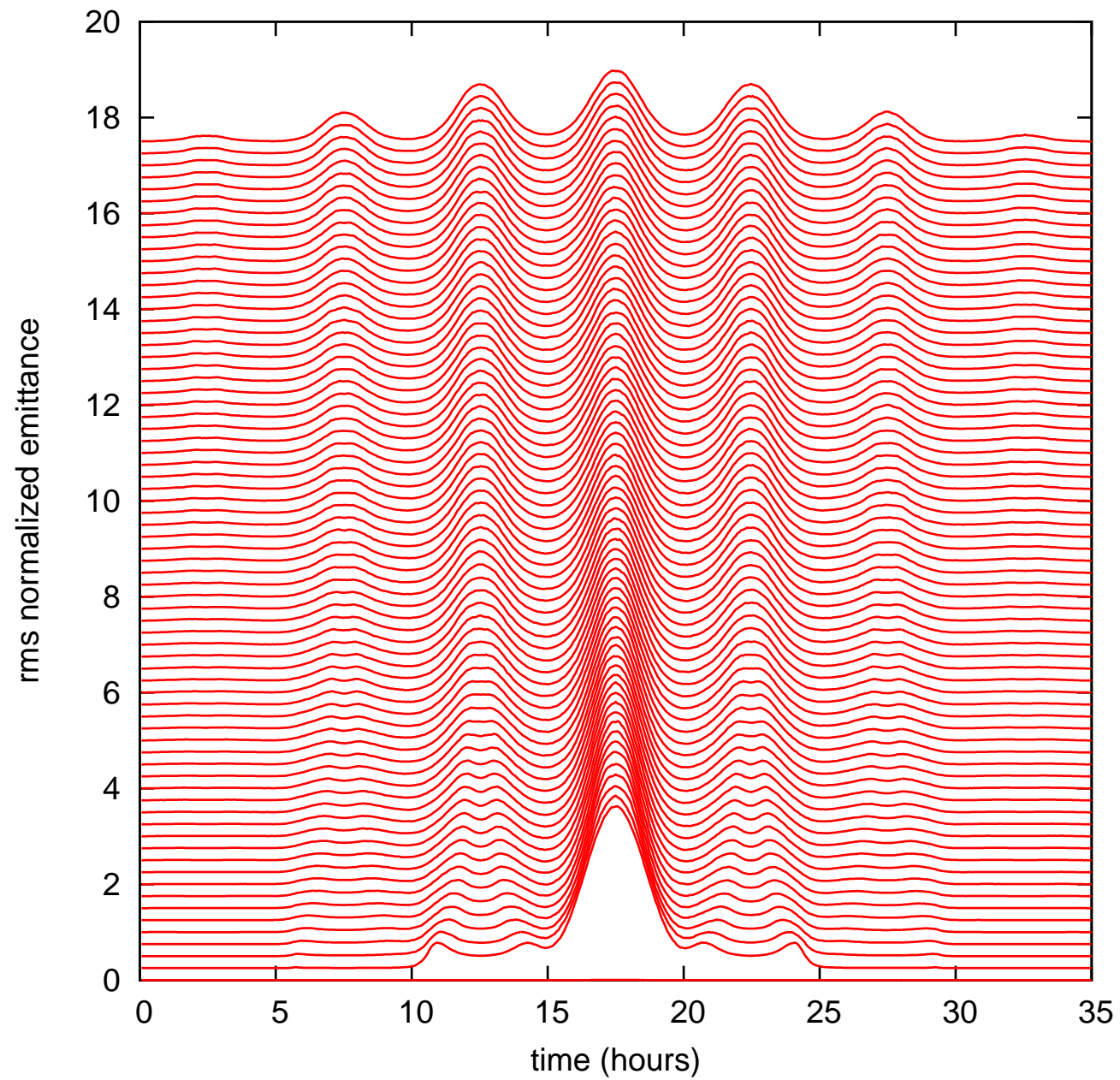


fibsx=0.5, Zs=3.e7,Zy=1.0e9

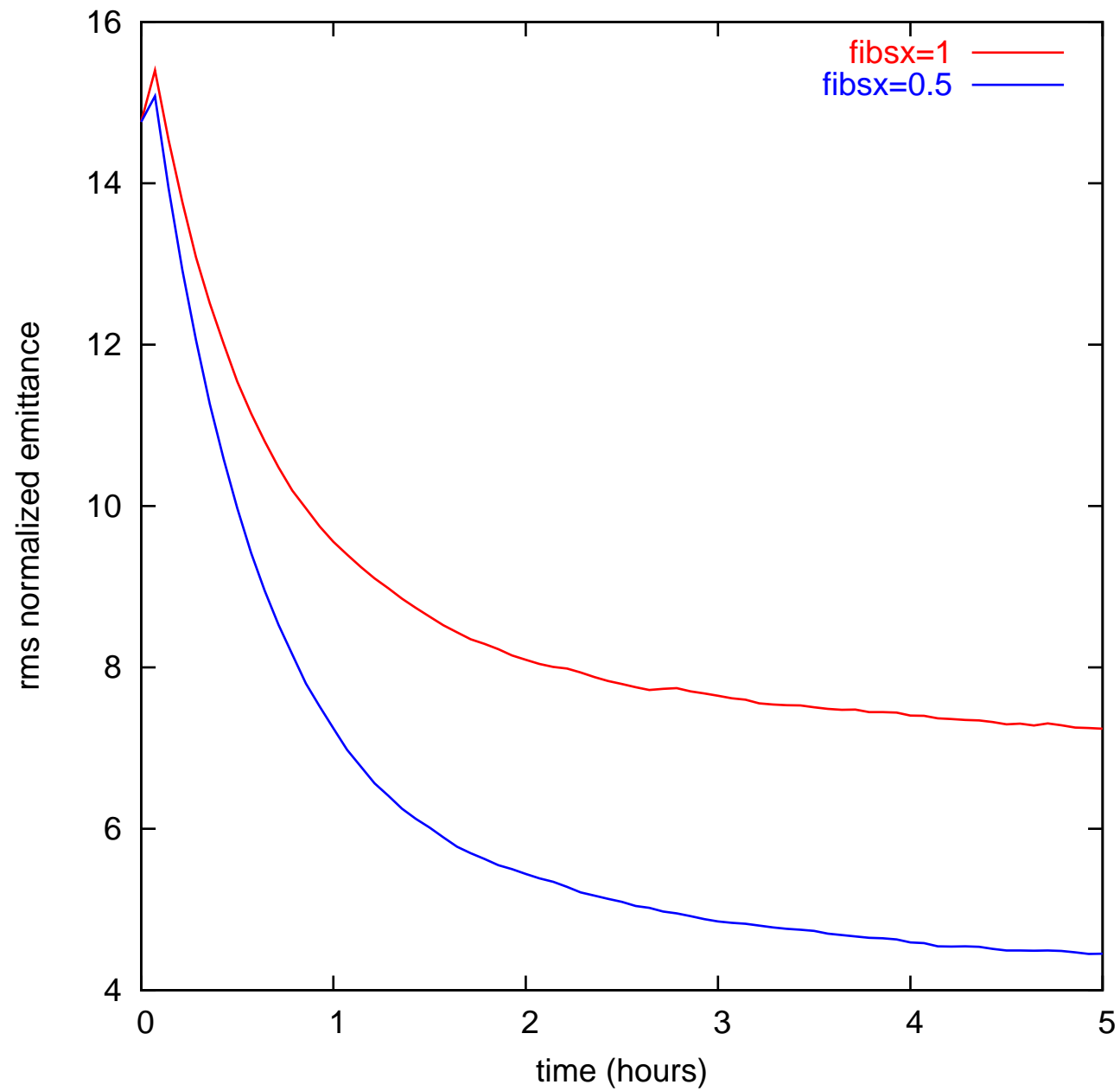




fibsx=0.5, Zs=3.e7,Zy=0.5e9



$Z_s=3.e7, Z_y=1.0e9$



$Z_s=3.e7, Z_y=0.5e9$

